

THE LACK OF COMPACTNESS IN THE SOBOLEV-STRICHARTZ INEQUALITIES

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ABSTRACT. We provide a general method to decompose any bounded sequence in \dot{H}^s into linear dispersive profiles generated by an abstract propagator, with a rest which is small in the associated Strichartz norms. The argument is quite different from the one proposed by Bahouri-Gérard and Keraani in the cases of the wave and Schrödinger equations, and is adaptable to a large class of propagators, including those which are matrix-valued.

1. INTRODUCTION

In the recent years, the research on nonlinear PDE's produced a relevant increment of strategies and techniques finalized to a complete understanding of some critical differential models. As a starting example, motivated by the interest on the Yamabe problem, some pioneer results were obtained by Aubin and Talenti in [1, 26], giving answers to some natural questions related to the criticality of the Sobolev embedding $\dot{H}^s(\mathbb{R}^d) \subset L^{p(s)}(\mathbb{R}^d)$, with $p(s) = 2d/(d-2s)$, and $0 < s < d/2$. Some years later, a great and well celebrated contribution to the theory of critical elliptic PDE's was given by Pierre Louis Lions, who introduced the concentration-compactness method, which immediately turned out to be a standard tool (see [19, 20, 21]). After the work by Lions, Solimini and Gérard in [10, 24] independently, and with different proofs, were able to describe in a precise way the lack of compactness of the Sobolev embedding $\dot{H}^s(\mathbb{R}^d) \subset L^{p(s)}(\mathbb{R}^d)$ (and also the version for Lorentz spaces, in [24]). Inspired to [10], Gallagher in [8], Bahouri and Gérard in [2] and Keraani in [15] proved analogous results related to the Sobolev-Strichartz estimates, respectively for the Navier-Stokes, the wave and the Schrödinger equation. As an example, we paste here the result proved by Keraani in [15]: the following standard notations

$$L_t^p L_x^q := L^p(\mathbb{R}; L^q(\mathbb{R}^d)), \quad L_t^p \dot{H}_x^s := L^p(\mathbb{R}; \dot{H}^s(\mathbb{R}^d)), \quad L_{t,x}^r := L_t^r L_x^r$$

will accompany the rest of the paper.

Theorem 1.1 (Keraani [15]). *Let $(\varphi_n)_{n \geq 0}$ be a bounded sequence in $\dot{H}^1(\mathbb{R}^3)$ and let $v_n(t, x) := e^{it\Delta} \varphi_n$. Then there exist a subsequence (v'_n) of (v_n) , a sequence $(\mathbf{h}^j)_{j \geq 1}$, $\mathbf{h}^j = (h_n^j)_{n \geq 0}$ for any $j \geq 1$ of scales, a sequence $(\mathbf{z}^j)_{j \geq 1} = (\mathbf{t}^j, \mathbf{x}^j)_{j \geq 1}$, with $\mathbf{z}^j = (t_n^j, x_n^j)_{n \geq 0}$ for any $j \geq 1$ of cores, and a sequence of functions $(U^j)_{j \geq 1}$ in $\dot{H}^1(\mathbb{R}^3)$ such that:*

$$(1.1) \quad \left| \frac{h_n^k}{h_n^j} \right| + \left| \frac{h_n^j}{h_n^k} \right| + \left| \frac{t_n^j - t_n^k}{(h_n^j)^2} \right| + \left| \frac{x_n^j - x_n^k}{h_n^j} \right| \rightarrow +\infty,$$

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as $n \rightarrow \infty$, for any $j \neq k$;

$$(1.2) \quad v'_n(t, x) = \sum_{j=1}^l \frac{1}{\sqrt{h_n^j}} e^{i\left(\frac{t-t_n^j}{(h_n^j)^2}\right)\Delta} U^j \left(\frac{x - x_n^j}{h_n^j} \right) + w_n^l(t, x),$$

for any $l \geq 1$, with

$$(1.3) \quad \limsup_{n \rightarrow \infty} \|w_n^l\|_{L_t^p L_x^q} \rightarrow 0,$$

as $l \rightarrow \infty$, for any (non-endpoint) H^1 -admissible couple (p, q) satisfying

$$\frac{2}{p} + \frac{3}{q} = \frac{3}{2} - 1, \quad 4 < p \leq \infty;$$

(1.4)

$$\int |\nabla_x v'_n(0, x)|^2 dx = \sum_{j=1}^l \int |\nabla U^j(x)|^2 dx + \int |\nabla_x w_n^l(0, x)|^2 dx + o(1), \quad \text{as } n \rightarrow \infty.$$

Almost in the same years of [2], [15], Kenig and Merle introduced in [13, 14] a new strategy to solve a large class of critical nonlinear Schrödinger and wave equations. The argument by Kenig and Merle is based on extrapolating, by contradiction, a single compactly behaving solution to the problem, which they call *critical element*, via concentration-compactness methods; then, the rigidity given by the algebra of the equation implies that such solution, with such compactness properties, cannot exist. The basic tool in capturing the critical element is given by a nonlinear version of Theorem 1.1 (in the case of Schrödinger, and the analogous in [2] for wave), which is in fact a consequence of the same result and the scattering properties of the nonlinear flow. Since the Kenig-Merle proof turns out to be adaptable to a large class of nonlinear dispersive equations, a lot of results appeared in the very last years in the same spirit of Theorem 1.1, for different propagators (see e.g. [3, 4, 5, 8, 9, 16, 17, 22, 23]. Among the previous list, we mention the papers by Merle-Vega [17], Begout-Vargas [3], Rogers-Vargas [23] and recently Ramos [22], in which Strichartz estimates at the lowest scales are treated, and some refinements are needed, in the style of the one which has been proved by Moyua-Vargas-Vega in [18]; in the cases of [2] and [15] the inequality (2.38) in the sequel, proved by Gérard in [10], plays the analog role of the Strichartz refinement.

Therefore, it would be appreciable to have a general result, in the same style of Theorem 1.1, which might hold for a large and unified class of dispersive propagators. On the other hand, as far as we can see, it is not clear if the strategy proposed in [2] and [15] might be adaptable, in total generality, to many problems, as for example the case of dispersive systems.

In view of the above considerations, the aim of this paper is to provide a new proof, which is quite different from the one proposed in [2] and [15], and which works for a large amount of dispersive propagators, including among the others the matrix-valued cases.

We are now ready to prepare the setting of our main theorem. In the following, we work with vector-valued functions $f = (f_1, \dots, f_N) : \mathbb{R}^d \rightarrow \mathbb{C}^N$, with the notation

$$(1.5) \quad \|f\|_{\dot{H}_x^s}^2 = \sum_{j=1}^N \|f_j\|_{\dot{H}_x^s}^2.$$

With the symbol $\mathcal{L} = \mathcal{L}(D)$ we denote an operator

$$\mathcal{L}(D) = \mathcal{F}^{-1}(\mathcal{L}(\xi)\mathcal{F}), \quad \mathcal{L}(\xi) = (\mathcal{L}_{ij}(\xi))_{i,j=1\dots N} : \mathbb{R}^d \rightarrow \mathcal{M}_{N \times N}(\mathbb{C}),$$

where \mathcal{F} is the standard Fourier transform, and the matrix $\mathcal{L}(\xi)$ is assumed to be hermitian; in the above setting, the dispersive character of the Cauchy problem

$$(1.6) \quad \begin{cases} i\partial_t u + \mathcal{L}(D)u = 0 \\ u(0, x) = f(x) \end{cases}$$

just depends on the geometrical properties of the graph of $\mathcal{L}(\xi)$, as it is well known. In addition, we make the following abstract assumptions:

(H1) there exists $0 < s < \frac{d}{2}$ such that the problem (1.6) is globally well-posed in \dot{H}_x^s , and the unique solution is given via the propagator $u(t, x) = e^{it\mathcal{L}(D)}f(x)$;

(H2) the flow $e^{it\mathcal{L}(D)}$ is unitary onto \dot{H}_x^s , i.e.

$$\|e^{it\mathcal{L}(D)}f\|_{\dot{H}_x^s} = \|f\|_{\dot{H}_x^s}, \quad \forall t \in \mathbb{R},$$

where s is the same as in (H1), and $\|\cdot\|_{\dot{H}_x^s}$ is defined in (1.5).

(H3) the symbol $\mathcal{L}(\xi) : \mathbb{R}^d \rightarrow \mathcal{M}_{N \times N}(\mathbb{C})$ is α -homogeneous, i.e., for all $\lambda > 0$,

$$\mathcal{L}(\lambda\xi) = \lambda^\alpha \mathcal{L}(\xi);$$

(H4) there exist $2 \leq p < q \leq \infty$ such that the following Strichartz estimate hold

$$\|e^{it\mathcal{L}(D)}f\|_{L_t^p L_x^q} \leq C\|f\|_{\dot{H}_x^s},$$

with the same s as in (H1) and some constant $C > 0$.

By homogeneity, the couple (p, q) in (H4) needs to satisfy the scaling condition

$$(1.7) \quad \frac{\alpha}{p} + \frac{d}{q} = \frac{d}{2} - s,$$

where s is given by (H1) and α is the one in (H3). Notice that, by the Sobolev embedding $\dot{H}_x^s \subset L_x^{\frac{2d}{d-2s}}$, for $0 < s < d/2$, and the \dot{H}_x^s -preservation

$$\|e^{it\mathcal{L}(D)}f\|_{L_t^\infty \dot{H}_x^s} = \|f\|_{\dot{H}_x^s}$$

(assumption (H2) above), we get

$$(1.8) \quad \|e^{it\mathcal{L}(D)}f\|_{L_t^\infty L_x^{\frac{2d}{d-2s}}} \leq C\|f\|_{\dot{H}_x^s}, \quad 0 < s < \frac{d}{2},$$

for some constant $C > 0$. Consequently, by interpolation with (1.8), an estimate as the one of assumption (H4) automatically holds for any s -admissible pair (\tilde{p}, \tilde{q}) , i.e. any (\tilde{p}, \tilde{q}) satisfying (1.7), with $\tilde{p} \geq p$. In particular, we have

$$(1.9) \quad \|e^{it\mathcal{L}(D)}f\|_{L_{t,x}^r} \leq C\|f\|_{\dot{H}_x^s}, \quad 0 < s < \frac{d}{2}, \quad r = \frac{2(\alpha + d)}{d - 2s}.$$

There are several examples of operators $\mathcal{L}(D)$ satisfying the previous assumptions, including the cases of Schrödinger, non-elliptic Schrödinger, wave and Dirac propagators, as we will show later during the introduction. We are now ready to state our main theorem.

Theorem 1.2. *Let $\mathbf{u} = (u_n)_{n \geq 0}$ be a bounded sequence in \dot{H}_x^s for $0 < s < \frac{d}{2}$. There exist a subsequence (u'_n) of (u_n) , a sequence $(\mathbf{h}^j)_{j \geq 1}$, $\mathbf{h}^j = (h_n^j)_{n \geq 0}$ of scales, for any $j \geq 1$, a sequence $(\mathbf{z}^j)_{j \geq 1} = (\mathbf{t}^j, \mathbf{x}^j)_{j \geq 1}$ of cores, with $\mathbf{z}^j = (t_n^j, x_n^j)_{n \geq 0}$ for any $j \geq 1$, and a sequence of functions $(U^j)_{j \geq 1}$ in \dot{H}_x^s such that:*

$$(1.10) \quad \left| \frac{h_n^m}{h_n^j} \right| + \left| \frac{h_n^j}{h_n^m} \right| + \left| \frac{t_n^j - t_n^m}{(h_n^j)^\alpha} \right| + \left| \frac{x_n^j - x_n^m}{h_n^j} \right| \rightarrow +\infty,$$

as $n \rightarrow \infty$, for any $j \neq m$;

$$(1.11) \quad u'_n(x) = \sum_{j=1}^J \frac{1}{(h_n^j)^{\frac{d}{2}-s}} e^{i \left(\frac{t_n^j}{(h_n^j)^\alpha} \right) \mathcal{L}(D)} U^j \left(\frac{x - x_n^j}{h_n^j} \right) + R_n^J(x),$$

where α is the one in (H3), for any $J \geq 1$, with

$$(1.12) \quad \limsup_{n \rightarrow \infty} \|e^{it\mathcal{L}(D)} R_n^J(x)\|_{L_t^{\tilde{p}} L_x^{\tilde{q}}} \rightarrow 0,$$

as $J \rightarrow \infty$, for any couple (\tilde{p}, \tilde{q}) satisfying the admissibility condition (1.7), with $p < \tilde{p} < \infty$, and p is the one given by (H4); for any $J \geq 1$ we have

$$(1.13) \quad \|u'_n(x)\|_{\dot{H}_x^s}^2 = \sum_{j=1}^J \|U^j(x)\|_{\dot{H}_x^s}^2 + \|R_n^J(x)\|_{\dot{H}_x^s}^2 + o(1), \quad \text{as } n \rightarrow \infty.$$

Remark 1.1. Notice that (1.11) is slightly different to (1.2); effectively, it is sufficient to act at both the sides of (1.11) with the propagator $e^{it\mathcal{L}(D)}$, to obtain the analogous of (1.2). In fact, we prefer to write (1.11) in this form, because it respect the *stationary* character of our proof. As it will be clear in the sequel, the main difference with the argument in [2], [15] is that, at each step of the recurrence argument which permits to extract the final sequence u'_n , we work on fixed sequences of times; arguing in this way, all the construction can be performed exactly as in the stationary theorem by Gérard in [10]. This idea is suggested by the argument which has been introduced in [7], to prove the existence of maximizers for Sobolev-Strichartz inequalities.

Remark 1.2. Theorem 1.2 implies Theorem 1.1, in the special cases $\mathcal{L}(D) = \Delta$, $s = 1$, $d = 3$, apart from (1.12); indeed, the case $p = \infty$ is missing in (1.12). We do not find possible to obtain the decay of the $L_t^\infty L_x^{p(s)}$ in total generality; on the other hand, it is possible to prove it case by case, using each time the specific properties of $\mathcal{L}(\xi)$. By the way, we stress that the decay of $L_{t,x}^r$ -norm in (1.12), when r is the one in (1.9), is typically the only information which is needed in the nonlinear applications.

We now pass to give some examples of applications of the main theorem to other types of propagators.

Example 1.1 (Wave propagator). The Strichartz estimates for the wave propagator $e^{it|D|}$ (see [11], [12]), in dimension $d \geq 2$, are the following:

$$(1.14) \quad \left\| e^{it|D|} f \right\|_{L_t^p L_x^q} \leq C \|f\|_{\dot{H}_x^{\frac{1}{p} - \frac{1}{q} + \frac{1}{2}}},$$

under the admissibility condition

$$(1.15) \quad \frac{2}{p} + \frac{d-1}{q} = \frac{d-1}{2}, \quad p \geq 2, \quad (p, q) \neq (2, \infty).$$

The gap of derivatives $\frac{1}{p} - \frac{1}{q} + \frac{1}{2} \geq 0$ is null only in the case of the energy estimate $(p, q) = (\infty, 2)$. In particular,

$$(1.16) \quad \left\| e^{it|D|} f \right\|_{L_{t,x}^{\frac{2(d+1)}{d-1}}} \leq C \|f\|_{\dot{H}_x^{\frac{1}{2}}} \quad d \geq 2,$$

which is in fact the original estimate proved by Strichartz in [25]. More generally, by Sobolev embedding one also obtains that

$$(1.17) \quad \left\| e^{it|D|} f \right\|_{L_{t,x}^{\frac{2(d+1)}{d-1-2\sigma}}} \leq C \|f\|_{\dot{H}_x^{\frac{1}{2}+\sigma}}, \quad 0 \leq \sigma < \frac{d-1}{2}, \quad d \geq 2.$$

Theorem 1.2 applies in this case, for any dimension $d \geq 2$, and $0 < \sigma < \frac{d-1}{2}$; notice that the case $\sigma = 0$ is not included, since in this case assumption (H4) fails. The case $\sigma = 0$ has been recently treated and solved by Ramos in [22].

Example 1.2 (Dirac propagator). In dimension $d = 3$, the massless Dirac operator is given by

$$\mathcal{D} := \frac{1}{i} \sum_{j=1}^3 \alpha_j \partial_j.$$

Here $\alpha_1, \alpha_2, \alpha_3 \in \mathcal{M}_4(\mathbb{C})$ are the so called *Dirac matrices*, which are 4×4 -hermitian matrices, $\alpha_j^t = \overline{\alpha_j}$, $j = 1, 2, 3$, with the explicit form

$$\alpha_1 = \begin{pmatrix} 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \\ 0 & 1 & 0 & 0 \\ 1 & 0 & 0 & 0 \end{pmatrix}, \quad \alpha_2 = \begin{pmatrix} 0 & 0 & 0 & -i \\ 0 & 0 & i & 0 \\ 0 & -i & 0 & 0 \\ i & 0 & 0 & 0 \end{pmatrix}, \quad \alpha_3 = \begin{pmatrix} 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & -1 \\ 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \end{pmatrix};$$

equivalently, $\alpha_j = \begin{pmatrix} 0 & \sigma_j \\ \sigma_j & 0 \end{pmatrix}$, where σ_j is the j^{th} 2×2 -Pauli matrix, $j = 1, 2, 3$.

Since $\mathcal{D}^2 = -\Delta_{4 \times 4}$, the Strichartz estimates for the massless Dirac operator are the same as for the 3D wave equation (see [6]):

$$(1.18) \quad \|e^{it\mathcal{D}} f\|_{L_t^p L_x^q} \leq C \|f\|_{\dot{H}_x^{\frac{1}{p} - \frac{1}{q} + \frac{1}{2}}}.$$

Here, the admissibility condition reads as follows:

$$(1.19) \quad \frac{2}{p} + \frac{2}{q} = 1, \quad p > 2.$$

In particular we have

$$(1.20) \quad \|e^{it\mathcal{D}} f\|_{L_{t,x}^{\frac{8}{2-2\sigma}}} \leq C \|f\|_{\dot{H}_x^{\frac{1}{2}+\sigma}}, \quad 0 \leq \sigma < 1.$$

Also in this case Theorem 1.2 applies; moreover, the statement also includes the cases of more general dispersive systems. At our knowledge, this is not a known fact; indeed, it is unclear if it might be possible to prove the same using the arguments by Bahouri-Gérard and Keraani in [2], [15].

Example 1.3 (Non-elliptic Schrödinger propagators). Let us consider the Schrödinger operator $L := \sum_{j=1}^m \partial_j^2 - \sum_{j=m+1}^d \partial_j^2$, for $d \geq 2$ and $1 \leq m < d$. The Strichartz estimates are the same as for the Schrödinger propagator, namely

$$(1.21) \quad \|e^{itL}f\|_{L_t^p L_x^q} \leq C\|f\|_{\dot{H}_x^s},$$

with the admissibility condition

$$(1.22) \quad \frac{2}{p} + \frac{d}{q} = \frac{d}{2} - s, \quad p \geq 2, \quad (p, q) \neq (2, \infty).$$

In particular, one has

$$(1.23) \quad \|e^{itL}f\|_{L_{t,x}^{\frac{2(d+2)}{d-2s}}} \leq C\|f\|_{\dot{H}_x^s},$$

for any $0 \leq s < \frac{d}{2}$. The only case in which at our knowledge has been treated is $d = 2, s = 0$. Indeed, the result by Rogers and Vargas in [23] contains all the ingredients which are necessary to prove the profile decomposition for bounded sequences in L^2 , with respect to the propagator e^{itL} , in dimension $d = 2$. It is a matter of fact that Theorem 1.2 applies for any $d \geq 2$ and $0 < s < \frac{d}{2}$, but we remark that it cannot include the case $s = 0$. In addition, our argument is rather simple and does not involve any Fourier properties of the propagator, since it just use the fact that $s > 0$.

The rest of the paper is devoted to the proof of Theorem 1.2.

2. PROOF OF THEOREM 1.2

Let us start with some preliminary definitions, introduced in [10].

Definition 2.1. Let $\mathbf{f} = (f_n)_{n \geq 1}$ be a bounded sequence in $L^2(\mathbb{R}^d)$ and $\mathbf{h} = (h_n)_{n \geq 1}, \tilde{\mathbf{h}} = (\tilde{h}_n)_{n \geq 1} \subset \mathbb{R}$ two scales. We say that:

- \mathbf{f} is \mathbf{h} -oscillatory if

$$(2.1) \quad \limsup_{n \rightarrow \infty} \left(\int_{h_n|\xi| \leq \frac{1}{R}} |\widehat{f_n}(\xi)|^2 d\xi + \left(\int_{h_n|\xi| \geq R} |\widehat{f_n}(\xi)|^2 d\xi \right) \right) \rightarrow 0 \quad \text{as } R \rightarrow \infty;$$

- \mathbf{f} is \mathbf{h} -singular if, for every $b > a > 0$, we have

$$(2.2) \quad \lim_{n \rightarrow \infty} \int_{a \leq h_n|\xi| \leq b} |\widehat{f_n}(\xi)|^2 d\xi = 0 \quad \text{as } R \rightarrow \infty;$$

- \mathbf{h} and $\tilde{\mathbf{h}}$ are orthogonal if

$$(2.3) \quad \lim_{n \rightarrow \infty} \left(\frac{h_n}{\tilde{h}_n} + \frac{\tilde{h}_n}{h_n} \right) = 0.$$

The following proposition, proved in [10], permits to reduce the matters to prove Theorem 1.2 in the case of $\mathbf{1}$ -oscillating sequences.

Proposition 2.1 (Gérard [10]). *Let $\mathbf{f} = (f_n)_{n \geq 0}$ be a bounded sequence in $L^2(\mathbb{R}^d)$. There exist a subsequence \mathbf{f}' of \mathbf{f} , a family $(\mathbf{h}^j)_{j \geq 1}$ of pairwise orthogonal scales and a family $(\mathbf{g}^j)_{j \geq 1}$ of bounded sequences in $L^2(\mathbb{R}^d)$ such that:*

- \mathbf{g}^j is \mathbf{h}^j -oscillatory, for every j ;

- for every $J \geq 1$ and $x \in \mathbb{R}^d$,

$$(2.4) \quad f'_n(x) = \sum_{j=1}^J g_n^j(x) + R_n^J(x),$$

where $(R_n^J)_{n \geq 1}$ is \mathbf{h}^J -singular, for every $j = 1, \dots, J$ and

$$(2.5) \quad \limsup_{n \rightarrow \infty} \|R_n^J\|_{\dot{B}_{2,\infty}^0} \rightarrow 0 \quad \text{as } J \rightarrow \infty;$$

- for every $J \geq 1$,

$$(2.6) \quad \|f'_n\|_{L^2}^2 = \sum_{j=1}^J \|g_n^j\|_{L^2}^2 + \|R_n^J\|_{L^2}^2 + o(1), \quad \text{as } n \rightarrow \infty.$$

We are now ready to prove the following result, which is the core of the proof of Theorem 1.2.

Proposition 2.2. *Assume (H1)-(H2)-(H3)-(H4). Let $\mathbf{u} = (u_n) \subset \dot{H}_x^s$ be a 1-oscillatory, bounded sequence in \dot{H}_x^s with $0 < s < \frac{d}{2}$. There exist a subsequence $\mathbf{u}' = (u'_n)$ of \mathbf{u} , a family of cores $(\mathbf{z}^j)_{j \geq 1} = (t_n^j, x_n^j)_{n \geq 0, j \geq 1} \subset \mathbb{R} \times \mathbb{R}^d$, and a family of functions $(U^j)_{j \geq 1}$ in \dot{H}_x^s such that:*

- (i) for any $j \neq k$, we have

$$(2.7) \quad |t_n^j - t_n^k| + |x_n^j - x_n^k| \rightarrow \infty, \quad \text{as } n \rightarrow \infty;$$

- (ii) for all $J \geq 1$ and $x \in \mathbb{R}^d$,

$$(2.8) \quad u'_n(x) = \sum_{j=1}^J e^{it_n^j \mathcal{L}(D)} U^j(x - x_n^j) + R_n^J(x),$$

with

$$(2.9) \quad \limsup_{n \rightarrow \infty} \|e^{it \mathcal{L}(D)} R_n^J(x)\|_{L_t^{\tilde{p}} L_x^{\tilde{q}}} \rightarrow 0, \quad \text{as } J \rightarrow \infty,$$

for any s -admissible pair (\tilde{p}, \tilde{q}) , with $\tilde{p} > p$, and p given by (H4). In addition,

$$(2.10) \quad \|u'_n(x)\|_{\dot{H}_x^s}^2 = \sum_{j=1}^J \|U^j(x)\|_{\dot{H}_x^s}^2 + \|R_n^J(x)\|_{\dot{H}_x^s}^2 + o(1), \quad \text{as } n \rightarrow \infty.$$

Proof. Let us introduce the notation $S := L_t^\infty L_x^{\frac{2d}{d-2s}}$, and recall (1.8). The proof is based on a construction by recurrence, which is quite different from the one used in [2], [15].

Assume that $\liminf_{n \rightarrow \infty} \|e^{-it \mathcal{L}(D)} u_n(x)\|_S = 0$, then (2.8) is satisfied provided that u'_n is a subsequence of u_n such that $\lim_{n \rightarrow \infty} \|e^{-it \mathcal{L}(D)} u'_n(x)\|_S = 0$, $J = 0$, $U^0 \equiv 0$ and $R_n^0 \equiv u'_n$; moreover in this case (2.7) and (2.10) are trivially satisfied and the decay of the interpolated Strichartz norms (2.9) follows by interpolation between the decay of the S -norm and the $L_t^p L_x^q$ a priori bound given by assumption (H4). Therefore we shall assume that

$$\|e^{-it \mathcal{L}(D)} u_n(x)\|_S \geq 2\delta_0,$$

for any $n \geq 0$ and some $\delta_0 > 0$. By the definition of S , there exists a sequence of times $(t_n^1)_{n \geq 0} \subset \mathbb{R}$ such that

$$(2.11) \quad w_n^1(x) := e^{-it_n^1 \mathcal{L}(D)} u_n(x), \quad \|w_n^1(x)\|_{L_x^{\frac{2d}{d-2s}}} \geq \frac{1}{2} \left\| e^{-it \mathcal{L}(D)} u_n(x) \right\|_S \geq \delta_0.$$

Arguing as Gérard in [10], we now denote by $\mathcal{P}(\mathbf{w}^1)$ the set of all the possible weak limits in \dot{H}_x^s of all the possible subsequences of (w_n^1) with all their possible translations; moreover, let

$$(2.12) \quad \gamma(\mathbf{w}^1) := \sup \left\{ \|\psi\|_{\dot{H}_x^s} : \psi \in \mathcal{P}(\mathbf{w}^1) \right\}.$$

As a consequence, there exist a sequence of centers $(x_n^1)_{n \geq 0} \subset \mathbb{R}^d$ and a subsequence of u_n (that we still denote u_n) such that

$$(2.13) \quad e^{-it_n^1 \mathcal{L}(D)} u_n(x + x_n^1) = w_n^1(x + x_n^1) \rightharpoonup U^1(x) \quad \text{weakly in } \dot{H}_x^s,$$

as $n \rightarrow \infty$, where

$$(2.14) \quad \gamma(\mathbf{w}^1) \leq 2 \|U^1\|_{\dot{H}_x^s}.$$

Next we introduce $R_n^1(x)$ as follows

$$(2.15) \quad u_n(x) = e^{it_n^1 \mathcal{L}(D)} U^1(x - x_n^1) + R_n^1(x),$$

and by (2.13) we get

$$(2.16) \quad e^{-it_n^1 \mathcal{L}(D)} R_n^1(x + x_n^1) \rightharpoonup 0 \quad \text{weakly in } \dot{H}_x^s,$$

as $n \rightarrow \infty$. Next notice that by combining (2.13) with (2.15) and by recalling the definition of weak limit we deduce

$$(2.17) \quad \left\| e^{-it_n^1 \mathcal{L}(D)} R_n^1(x + x_n^1) \right\|_{\dot{H}_x^s}^2 = \left\| e^{-it_n^1 \mathcal{L}(D)} u_n(x + x_n^1) \right\|_{\dot{H}_x^s}^2 - \|U^1(x)\|_{\dot{H}_x^s}^2 + o(1),$$

and by the \dot{H}_x^s -preservation (H2) implies

$$(2.18) \quad \|R_n^1(x)\|_{\dot{H}_x^s}^2 = \|u_n(x)\|_{\dot{H}_x^s}^2 - \|U^1(x)\|_{\dot{H}_x^s}^2 + o(1).$$

Next assume that

$$(2.19) \quad \liminf_{n \rightarrow \infty} \left\| e^{it \mathcal{L}(D)} R_n^1(x) \right\|_S = 0$$

then (2.8), (2.10) follow by (2.15), (2.18), and (2.9) follows by interpolation between the S -norm (that goes to zero on a suitable subsequence due to (2.19)) with the Strichartz norm given by assumption (H4). Therefore, up to choose a subsequence, we can assume as before that

$$\left\| e^{it \mathcal{L}(D)} R_n^1(x) \right\|_S \geq 2\delta_1,$$

for any $n \geq 0$ and some $\delta_1 > 0$. As a consequence, there exists a sequence of times $(t_n^2)_{n \geq 0} \subset \mathbb{R}$ such that

$$(2.20) \quad w_n^2(x) := e^{-it_n^2 \mathcal{L}(D)} R_n^1(x), \quad \|w_n^2(x)\|_{L_x^{\frac{2d}{d-2s}}} \geq \frac{1}{2} \left\| e^{-it \mathcal{L}(D)} R_n^1(x) \right\|_S \geq \delta_1.$$

Define as above

$$(2.21) \quad \gamma(\mathbf{w}^2) := \sup \left\{ \|\psi\|_{\dot{H}_x^s} : \psi \in \mathcal{P}(\mathbf{w}^2) \right\}.$$

To this we associate a new sequence $(x_n^2)_{n \geq 0} \subset \mathbb{R}^d$ and a subsequence of R_n^1 , which we still call R_n^1 , such that

$$(2.22) \quad e^{-it_n^2 \mathcal{L}(D)} R_n^1(x + x_n^2) = w_n^2(x + x_n^2) \rightharpoonup U^2(x) \quad \text{weakly in } \dot{H}_x^s,$$

as $n \rightarrow \infty$; moreover we can assume that

$$(2.23) \quad \gamma(\mathbf{w}^2) \leq 2 \|U^2\|_{\dot{H}_x^s}.$$

Next we introduce $R_n^2(x)$ as follows:

$$(2.24) \quad R_n^1(x) = e^{it_n^2 \mathcal{L}(D)} U^2(x - x_n^2) + R_n^2(x),$$

By (2.22) we conclude that

$$(2.25) \quad e^{-it_n^2 \mathcal{L}(D)} R_n^2(x + x_n^2) \rightharpoonup 0 \quad \text{weakly in } \dot{H}^s,$$

Moreover arguing as in (2.18) we get

$$(2.26) \quad \|R_n^2(x)\|_{\dot{H}_x^s}^2 = \|R_n^1(x)\|_{\dot{H}_x^s}^2 - \|U^2(x)\|_{\dot{H}_x^s}^2 + o(1).$$

By combining (2.15) and (2.24) we obtain

$$(2.27) \quad u_n(x) = e^{it_n^1 \mathcal{L}(D)} U^1(x - x_n^1) + e^{it_n^2 \mathcal{L}(D)} U^2(x - x_n^2) + R_n^2(x),$$

and by combining (2.26) with (2.18) we get

$$(2.28) \quad \|R_n^2(x)\|_{\dot{H}_x^s}^2 = \|u_n(x)\|_{\dot{H}_x^s}^2 - \|U^1(x)\|_{\dot{H}_x^s}^2 - \|U^2(x)\|_{\dot{H}_x^s}^2 + o(1).$$

The computations above describe an iterative procedure which at any step $j = 0, 1, \dots$ permits to construct a (finite) family $U^1, \dots, U^j \in \dot{H}_x^s$, a family of cores $(t_n^1, x_n^1)_{n \geq 1}, \dots, (t_n^j, x_n^j)_{n \geq 0} \in \mathbb{R} \times \mathbb{R}^d$, and a sequence $R_n^j(x) \in \dot{H}_x^s$ such that (up to subsequence) u_n can be written as

$$(2.29) \quad u_n(x) = e^{it_n^1 \mathcal{L}(D)} U^1(x - x_n^1) + \dots + e^{it_n^j \mathcal{L}(D)} U^j(x - x_n^j) + R_n^j(x),$$

with the following extra properties:

$$(2.30) \quad \|e^{-it_n^j \mathcal{L}(D)} R_n^{j-1}(x)\|_{L_x^{\frac{2d}{d-2s}}} \geq \frac{1}{2} \|e^{-it \mathcal{L}(D)} R_n^{j-1}(x)\|_S \geq \delta_{j-1} > 0;$$

$$(2.31) \quad e^{-it_n^j \mathcal{L}(D)} R_n^{j-1}(x + x_n^j) \rightharpoonup U^j(x) \quad \text{weakly in } \dot{H}_x^s;$$

$$(2.32) \quad \gamma(\mathbf{w}^j) \leq 2 \|U_j\|_{\dot{H}_x^s}$$

(where the sequence (w_n^j) is defined by $e^{-it_n^j \mathcal{L}(D)} R_n^{j-1}(x)$) and $\gamma(\mathbf{w}^j)$ is defined according to (2.12));

$$(2.33) \quad \|R_n^j(x)\|_{\dot{H}_x^s}^2 = \|u_n(x)\|_{\dot{H}_x^s}^2 - \|U^1(x)\|_{\dot{H}_x^s}^2 - \dots - \|U^j(x)\|_{\dot{H}_x^s}^2 + o(1);$$

$$(2.34) \quad e^{-it_n^j \mathcal{L}(D)} R_n^j(x + x_n^j) \rightharpoonup 0, \text{ weakly in } \dot{H}_x^s.$$

Notice that (2.29) and (2.33) prove (2.8) and (2.10). Our goal is now to prove (2.9); it is sufficient to prove that, for any $\epsilon > 0$, there exists $J = J(\epsilon) \in \mathbb{N}$ such that, for any $n \in \mathbb{N}$, and for any $j \geq J(\epsilon)$ we have

$$(2.35) \quad \limsup_{n \rightarrow \infty} \|e^{it \mathcal{L}(D)} R_n^j(x)\|_S < \epsilon.$$

Fix $\epsilon > 0$; first observe that, by (2.33) the sum $\sum_{j \geq 1} \|U^j\|_{\dot{H}_x^s}^2$ has to converge, then there exists $J = J(\epsilon)$ such that, for any $j \geq J(\epsilon)$,

$$(2.36) \quad \gamma(\mathbf{w}^j) \leq 2 \|U^j\|_{\dot{H}_x^s} < 2\epsilon$$

where we have used (2.32). In order to conclude (2.35) it is sufficient, by (2.30), to prove that

$$(2.37) \quad \|e^{-it_n^j \mathcal{L}(D)} R_n^{j-1}(x)\|_{L_x^{\frac{2d}{d-2s}}} := \|w_n^j(x)\|_{L_x^{\frac{2d}{d-2s}}} < C\epsilon,$$

for any $j \geq J(\epsilon)$, and some constant $C > 0$. This is an immediate consequence of the inequality

$$(2.38) \quad \limsup_{n \rightarrow \infty} \|w_n^j\|_{L_x^{p(s)}} \leq C \limsup_{n \rightarrow \infty} \|w_n^j\|_{\dot{H}_x^s}^{\frac{2}{p(s)}} \gamma(\mathbf{w}^j)^{1 - \frac{2}{p(s)}},$$

with $p(s) = 2d/(d-2s)$. The previous estimate has been proved by Gérard (see [10], estimate (4.19)).

In order to complete the proof, we need to show the orthogonality of the cores (2.7). Let us first prove it in the case $k = j+1$. Notice that by (2.31) we have

$$(2.39) \quad e^{-it_n^{j+1} \mathcal{L}(D)} R_n^j(x + x_n^{j+1}) \rightharpoonup U^{j+1}(x) \quad \text{weakly in } \dot{H}_x^s$$

which is equivalent to

$$e^{-i((t_n^{j+1} - t_n^j) + t_n^j) \mathcal{L}(D)} R_n^j(x + (x_n^{j+1} - x_n^j) + x_n^j) \rightharpoonup U^{j+1}(x).$$

Next assume by the absurd that the cores (\mathbf{z}^j) and (\mathbf{z}^{j+1}) do not satisfy (2.7), then up to subsequence we can assume $t_n^{j+1} - t_n^j \rightarrow \bar{t}$ and $x_n^{j+1} - x_n^j \rightarrow \bar{x}$, which in turn implies

$$e^{-i(\bar{t} + t_n^j) \mathcal{L}(D)} R_n^j(x + \bar{x} + x_n^j) \rightharpoonup U^{j+1}(x).$$

Notice that this last fact is in contradiction with (2.34).

Next we assume that there exist a couple (k, j) such that $k < j-1$ and for which the orthogonality (2.7) for the cores $(\mathbf{z}^j), (\mathbf{z}^k)$ is false (the case $k = j-1$ has been treated above). Moreover we can suppose that the orthogonality relation is satisfied for the cores (\mathbf{z}^{k+r}) and (\mathbf{z}^j) for any $r = 1, \dots, j-k-1$. In fact it is sufficient to choose k as

$$\sup\{h < j-1 \mid (\mathbf{z}^h) \text{ is not orthogonal to } (\mathbf{z}^j)\}.$$

Next notice that

$$R_n^k(x) = \sum_{h=k+1}^{j-1} e^{it_n^h \mathcal{L}(D)} U_h(x - x_n^h) + R_n^{j-1}(x)$$

(to prove this fact apply (2.29) twice: first up to the reminder R_n^k and after up to the reminder R_n^{j-1} , and subtract the two identities). As a consequence we get

$$(2.40) \quad \begin{aligned} & e^{-it_n^j \mathcal{L}(D)} R_n^k(x + x_n^j) \\ &= \sum_{h=k+1}^{j-1} e^{i(t_n^h - t_n^j) \mathcal{L}(D)} U_h(x + x_n^j - x_n^h) + e^{-it_n^j \mathcal{L}(D)} R_n^{j-1}(x + x_n^j). \end{aligned}$$

Next notice that by the orthogonality of (\mathbf{z}^{k+r}) and (\mathbf{z}^j) for $r = 1, \dots, j-k-1$ we get

$$(2.41) \quad e^{i(t_n^h - t_n^j) \mathcal{L}(D)} U_h(x + x_n^j - x_n^h) \rightharpoonup 0$$

for every $h = k + 1, \dots, j - 1$ (here we use Lemma A.1). On the other hand we have the following identity

$$e^{-it_n^j \mathcal{L}(D)} R_n^k(x + x_n^j) = e^{-i(t_n^j - t_n^k + t_n^k) \mathcal{L}(D)} R_n^k(x + x_n^j - x_n^k + x_n^k)$$

and since we are assuming that (\mathbf{z}^j) and (\mathbf{z}^k) are not orthogonal, then by compactness we can assume that $x_n^j - x_n^k \rightarrow \bar{x}$ and $t_n^j - t_n^k \rightarrow \bar{t}$. In particular we get

$$e^{-it_n^j \mathcal{L}(D)} R_n^k(x + x_n^j) - e^{-i(\bar{t} + t_n^k) \mathcal{L}(D)} R_n^k(x + \bar{x} + x_n^k) \rightarrow 0$$

and since by (2.34) we have $e^{-it_n^k \mathcal{L}(D)} R_n^k(x + x_n^k) \rightharpoonup 0$, then necessarily also the l.h.s. in (2.40) converges weakly to zero. By combining this fact with (2.41) we deduce that

$$e^{-it_n^j \mathcal{L}(D)} R_n^{j-1}(x + x_n^j) \rightharpoonup 0$$

and it is in contradiction with (2.31). \square

Having in mind Proposition 2.1, to complete the proof of Theorem 1.2 it is now sufficient to follow exactly the arguments given by Keraani in [15]. One should only be careful at the moment of proving (1.12); indeed, notice that in Lemma 2.7 in [15] it is used the fact that $r = 10$ is an integer number. On the other hand, the reader should easily notice that this is not a relevant fact, and the proof can be easily performed in the general case in which r is given by (1.9). Once the decay of the $L_{t,x}^r$ of the rest is proved, the decay of the norms in (1.12) follows by interpolation. We omit here further details.

APPENDIX A.

We devote this small appendix to prove a general result, Lemma A.1, which has been implicitly used during the proof of Proposition 2.2. Let us start with the following proposition.

Proposition A.1. *Assume that*

$$(A.1) \quad \left\| e^{it\mathcal{L}(D)} f \right\|_{L_t^p L_x^q} \leq C \|f\|_{\dot{H}_x^s},$$

for some $p, q \geq 1$, $p \neq \infty$ and some $C > 0$. Then $\|e^{it\mathcal{L}(D)} f\|_{L_x^q} \rightarrow 0$, as $t \rightarrow \infty$, for any $f \in C_0^\infty(\mathbb{R}^d)$.

Proof. Let $M \gg 1$ such that $H_x^M \subset L_x^q$, and let $f \in H_x^M$. Then, by continuity in time, for every $\epsilon > 0$ there exists $\bar{t} = \bar{t}(\epsilon, f) > 0$ such that

$$(A.2) \quad \|e^{it\mathcal{L}(D)} f - f\|_{H_x^M} \leq \epsilon, \quad \forall |t| < \bar{t}.$$

Now assume by the absurd that for a sequence $t_n \rightarrow \infty$ we have

$$(A.3) \quad \inf_n \left\| e^{it_n \mathcal{L}(D)} f \right\|_{L_x^q} = \delta > 0.$$

As a consequence, by combining the Sobolev embedding $H_x^M \subset L_x^q$ with (A.2) and the fact that $e^{it\mathcal{L}(D)}$ is an isometry on H_x^M , we have

$$\begin{aligned} \|e^{i(t_n+h)\mathcal{L}(D)} f - e^{it_n \mathcal{L}(D)} f\|_{L_x^q} &\leq C \|e^{i(t_n+h)\mathcal{L}(D)} f - e^{it_n \mathcal{L}(D)} f\|_{H_x^M} \\ &= C \|e^{ih\mathcal{L}(D)} f - f\|_{H_x^M} \leq \frac{\delta}{2}, \end{aligned}$$

provided that $|h| \leq h(\delta, f)$. Therefore we deduce by (A.3) that

$$\|e^{i(t_n+h)\mathcal{L}(D)}f\|_{L_x^q} \geq \frac{\delta}{2},$$

for any $n \in \mathbb{N}$ and $|h| \leq \bar{h}$. The last estimate is in contradiction with (A.1) since it does not allow global summability in time. \square

We can now prove the main result of the appendix.

Lemma A.1. *Assume that*

$$(A.4) \quad \left\| e^{it\mathcal{L}(D)}f \right\|_{L_t^p L_x^q} \leq C \|f\|_{\dot{H}_x^s},$$

for some $p, q \geq 1$ and some $C > 0$. Let $f \in \dot{H}_x^s$ and $\max\{|t_n|, |x_n|\} \rightarrow \infty$ then $e^{it_n\mathcal{L}(D)}f(x+x_n) \rightarrow 0$ in \dot{H}_x^s .

Proof. We need to consider two cases. The first possibility is that t_n is bounded; then necessarily x_n goes to ∞ and it is easy to conclude. In the case $t_n \rightarrow \infty$, the conclusion is now simple, by combining a density argument with Proposition A.1. \square

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